

Einführung in die Quantenoptik II

Sommersemester 2014

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Übungsaufgaben Blatt 4

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Problem 4.1 – Positive maps and the Kraus representation (10 points)

In the lecture, we have encountered dynamical (or completely positive = CP) maps. Recall that such a map $\Lambda : \rho \mapsto \Lambda(\rho)$

- is linear: $\Lambda(p\rho + q\rho') = p\Lambda(\rho) + q\Lambda(\rho')$ for real $p, q \geq 0$
- is positive and preserves the trace: $\langle \psi | \Lambda(\rho) | \psi \rangle \geq 0$ for all $|\psi\rangle$ and all ρ , $\text{tr } \Lambda(\rho) = \text{tr } \rho$.
- is completely positive: if Λ is extended to $\Lambda \otimes \mathbb{1}$ on a larger Hilbert space, then this map is positive.

(1) Show that unitary time evolution for a closed system, $\rho \mapsto U\rho U^\dagger$ with unitary U , is a CP map.

(2) Show that CP maps form a convex set: $p\Lambda + q\Lambda'$ is CP for real $0 \leq p, q \leq 1$, $p + q = 1$ if Λ and Λ' are CP.

(3) Conclude that the ‘Kraus operation’ defined as follows is a CP map:

$$\rho \mapsto \sum_k p_k U_k \rho U_k^\dagger, \quad \sum_k p_k = 1, \quad U_k \text{ unitary} \quad (4.1)$$

Problem 4.2 – Spectra – Fourier vs Laplace (10 points)

Correlation functions $\langle B(t)A(t') \rangle$ are nearly always calculated under stationary conditions, i.e., $\langle B(t+T)A(t'+T) \rangle = \langle B(t)A(t') \rangle$ for any time shift T .

(1) Show that stationary correlations only depend on the time difference $\tau = t - t'$.

(2) Show that if A and B are hermitean conjugates of each other, $A^\dagger = B$, then

$$\int_{-\infty}^{+\infty} d\tau e^{-i\omega\tau} \langle A^\dagger(t+\tau)A(t) \rangle = 2 \text{Re} \int_0^{+\infty} d\tau e^{-i\omega\tau} \langle A^\dagger(t+\tau)A(t) \rangle \quad (4.2)$$

so that one only needs ‘time-ordered (stationary) correlations’.

(3) Assume that the correlation function becomes a sum of damped exponentials (complex λ_k with $\text{Re } \lambda_k \geq 0$ are possible)

$$\langle A^\dagger(t+\tau)A(t) \rangle = \sum_k e^{-\lambda_k \tau} S_k \quad (4.3)$$

and show that the half-sided Fourier integral in Eq.(4.2) yields a spectrum which is a sum of Lorentzians. What happens when $\text{Re } \lambda_k \rightarrow 0$?

Problem 4.3 – Positive maps and transposition (10 points)

[Suggested reading for the interested student.] Transposition is an example of a positive, but not completely positive map. To see this, consider a two-level system and define the transposed density matrix in a given basis $|a\rangle, a = e, g$, by

$$\rho = \sum_{a,b} \rho_{ab} |a\rangle\langle b| \mapsto \rho^T = \sum_{a,b} \rho_{ab} |b\rangle\langle a| \quad (4.4)$$

(1) Show that the matrix representing ρ^T is the transpose of ρ ; its matrix elements are the complex conjugates of those of ρ : $(\rho^T)_{ab} = (\rho_{ab})^*$. Show that ρ^T is again a density operator.

(2) For a system of two qubits, take the basis $\{|ee\rangle, |ge\rangle, |eg\rangle, |gg\rangle\}$ and consider in this basis a two-qubit density matrix (greek capital ρ)

$$P = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \quad (4.5)$$

with 2×2 block matrices A, B, C, D . Check that $A^\dagger = A, D^\dagger = D, B^\dagger = C, \text{tr}(A + D) = 1$. Calculate the reduced density matrices $\text{tr}_{2,1}(P)$ for the first and the second qubit (the index at the trace means: ‘trace out this sub-system’)

$$\text{tr}_2(P) = A + D, \quad \text{tr}_1(P) = \begin{pmatrix} \text{tr } A & \text{tr } B \\ \text{tr } C & \text{tr } D \end{pmatrix} \quad (4.6)$$

(3) The transposed density matrix is of course

$$P^T = \begin{pmatrix} A^T & C^T \\ B^T & D^T \end{pmatrix} = \begin{pmatrix} A^* & B^* \\ C^* & D^* \end{pmatrix} \quad (4.7)$$

The *partial transpose* P^Γ (‘a half of a T ’, applied to the *second* system) is the operator defined by

$$P = \sum_{a,b,c,d} P_{ac,bd} |a c\rangle\langle b d| \mapsto P^\Gamma = \sum_{a,b,c,d} P_{ac,bd} |a d\rangle\langle b c| \quad (4.8)$$

Show that for two qubits, its matrix representation is

$$P^\Gamma = \begin{pmatrix} A & C \\ B & D \end{pmatrix} \quad (4.9)$$

‘This is intuitive’: the reduced state for the second qubit, $\text{tr}_1(P^\Gamma)$, is the transpose of $\text{tr}_1(P)$, while nothing changes when we trace out the second qubit. The

‘total transpose’ of this is the partial transpose on the first system, which we denote P^Γ

$$P^\Gamma = (P^\Gamma)^\Gamma = \begin{pmatrix} A^\Gamma & B^\Gamma \\ C^\Gamma & D^\Gamma \end{pmatrix} = \begin{pmatrix} A^* & C^* \\ B^* & D^* \end{pmatrix} \quad (4.10)$$

(4) Show that all these maps are involutions: apply them twice, and nothing has changed.

(5) Consider the state (this one is called ‘completely entangled’, a projector on $|ge\rangle + |eg\rangle$)

$$P = \begin{pmatrix} 0 & & & \\ & 1 & 1 & \\ & 1 & 1 & \\ & & & 0 \end{pmatrix} \quad (4.11)$$

where non-written matrix elements are zero, and show that its partial transposes $P^\Gamma = P^\Gamma$ are *not* positive.