

Einführung in die Quantenoptik II

Sommersemester 2016

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Übungsaufgaben Blatt 4

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Problem 4.1 – Phase diffusion model (8 points)

The process of ‘phase diffusion’ is a simple example of a Fokker-Planck equation. Consider an oscillatory signal whose phase $\phi = \phi(t)$ is drifting, following a random walk. (i) Write a sentence about the meaning of the distribution function $P(\phi, t)$. Choose and describe the initial distribution $P(\phi, 0)$ and justify that for large times,

$$\lim_{t \rightarrow \infty} P(\phi, t) = \frac{1}{2\pi} \quad (4.1)$$

(ii) The random walk dynamics of the phase can be described by the drift-diffusion equation

$$\frac{\partial P}{\partial t} - \omega_L \frac{\partial P}{\partial \phi} = D \frac{\partial^2 P}{\partial \phi^2} \quad (4.2)$$

Check that ω_L has the dimension of a frequency and D that of a diffusion coefficient. Transform into the ‘co-moving frame’

$$P(\phi, t) = \tilde{P}(\varphi, t), \quad \phi = \varphi - \omega_L t \quad (4.3)$$

and show that the solution can be written in the form

$$\tilde{P}(\varphi, t) = \sum_{n \in \mathbb{Z}} p_n e^{in\varphi - n^2 Dt} \quad (4.4)$$

(iii) Conclude that for a fixed-amplitude signal, the following average holds

$$\langle e^{i\phi(t)} \rangle = \langle e^{i\phi(0)} \rangle e^{-i\omega_L t} e^{-Dt} \quad (4.5)$$

Hint. Study first from Eq.(4.4) the phase distribution at $t = 0$.

Problem 4.2 – From master to Fokker-Planck equation (12 points)

In the lecture (and in many quantum optics books), we have found the following rules for the action of bosonic operators on coherent state projectors:

$$\begin{aligned} a|\alpha\rangle\langle\alpha| &= \alpha|\alpha\rangle\langle\alpha| \\ |\alpha\rangle\langle\alpha|a^\dagger &= \alpha^*|\alpha\rangle\langle\alpha| \\ a^\dagger|\alpha\rangle\langle\alpha| &= (\alpha^* + \partial_\alpha)|\alpha\rangle\langle\alpha| \\ |\alpha\rangle\langle\alpha|a &= (\alpha + \partial_{\alpha^*})|\alpha\rangle\langle\alpha| \end{aligned} \quad (4.6)$$

(i) To find an equation of motion for the P-function, an integration by parts (*partielle Integration*) is applied. Show that this gives, for example

$$\begin{aligned} a^\dagger a \rho &= \int d^2\alpha |\alpha\rangle\langle\alpha| (\alpha^* - \partial_\alpha) [\alpha P] \\ a^\dagger \rho a &= \int d^2\alpha |\alpha\rangle\langle\alpha| (\alpha^* - \partial_\alpha) (\alpha - \partial_{\alpha^*}) P \end{aligned} \quad (4.7)$$

(ii) Check that this ‘operator calculus’ is consistent with the law $(a^\dagger \rho) a = a^\dagger (\rho a)$ because the two differential operators $(\alpha^* - \partial_\alpha)$ and $(\alpha - \partial_{\alpha^*})$ commute. Check also that the first line in (4.7) and its counterpart for $aa^\dagger \rho$ are consistent with the commutator $[a, a^\dagger]$.

(iii) A reasonable master equation for a laser has the following form

$$\begin{aligned} \frac{d\rho}{dt} &= -i\omega_c [a^\dagger a, \rho] + \kappa a \rho a^\dagger - \frac{\kappa}{2} \{a^\dagger a, \rho\} \\ &\quad + G(a^\dagger a) a^\dagger \rho a - \frac{G(a^\dagger a)}{2} \{aa^\dagger, \rho\}, \end{aligned} \quad (4.8)$$

which is approximate because the nonlinear gain operator $G(a^\dagger a)$ should be handled more carefully. Analyze for example, whether the trace of ρ is conserved by the three terms ‘free evolution’, ‘cavity loss’, and ‘gain’.

(iv) Neglect first gain saturation and show that the equation resulting from (4.8) is (no guarantee for factors 1/2, please find the first term)

$$\begin{aligned} \frac{\partial}{\partial t} P(\alpha, \alpha^*) &= -i\omega_c \{ \dots \} P + \frac{1}{2} \{ \partial_\alpha (\kappa - G) \alpha + \partial_{\alpha^*} (\kappa - G) \alpha^* \} P \\ &\quad + \frac{G}{4} \partial_\alpha \partial_{\alpha^*} P, \end{aligned} \quad (4.9)$$

(The derivatives act on everything to their right, including P .) An approximate way to include gain saturation is to replace $G \mapsto G(|\alpha|^2) = G_0/(1 + B|\alpha|^2)$ and put it between the two derivatives in the last term. This term generates diffusion in phase-space. In the steady state, diffusion coefficient is, in order of magnitude, $D \sim G(\bar{n}) \sim \kappa$ where \bar{n} is the average photon number.