

# Introduction to Quantum Optics

Winter term 2006/07

Carsten Henkel/Martin Wilkens

## Problem set 3

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### Problem 3.1 – Q-function and coherent states (10 points)

We deal here with the quantum states for a single mode of the quantized field.

(1) Compute the Q-function for a thermal state of the mode.

(2) Compute the overlap between two coherent states  $|\alpha\rangle$  and  $|\beta\rangle$  and check that the Q-function for a coherent state is an isotropic gaussian in the complex  $\alpha$ -plane.

(3) The coherent states are not orthogonal, as the result of question (2) shows. Hence it is not surprising that they are ‘overcomplete’ in the following sense:

$$\int d^2\alpha |\alpha\rangle\langle\alpha| = \pi \mathbb{1} \quad (3.1)$$

where the unit operator is for example  $\mathbb{1} = \sum_{n=0}^{\infty} |n\rangle\langle n|$ . Prove this equation, using the standard integration measure in the complex plane:  $d^2\alpha = d(\operatorname{Re}\alpha) d(\operatorname{Im}\alpha)$ . The integral given in Problem 3.3 should be useful.

(4) Derive the following identity for an arbitrary density operator  $\rho$

$$\langle a^m a^{\dagger n} \rangle_{\rho} = \int \frac{d^2\alpha}{\pi} \alpha^m \alpha^{*n} Q_{\rho}(\alpha) \quad (3.2)$$

Speculate why this is usually interpreted as: *For averages in ‘anti-normal order’, the Q-function plays the role of a probability distribution.*

### Problem 3.2 – Casimir repulsion (10\* points)

Revisit the calculation of the Casimir energy that we did in the lecture for mixed ‘Dirichlet-von Neumann’ boundary conditions.

(2) Justify that the allowed wave vectors along the cavity axis are given by  $k_n L = \pi(n + \frac{1}{2})$  with  $n = 0, 1, 2, \dots$

(3) Perform the calculation of the Casimir energy in the same way as in the lecture. You will need the integral

$$\int_0^{\infty} \frac{dt t^3}{e^{2\pi t} + 1} = \frac{7}{1920} \quad (3.3)$$

(Bonus points.)

**Problem 3.3 – Dipole matrix elements (10 points)**

Consider the energy levels  $nl = 1s, 2s$  and  $2p$  of the Hydrogen atom. The coupling to the light field involves in the long-wavelength approximation the following matrix elements of the electron coordinate  $\mathbf{r}$  (measured relative to the nucleus):

$$-e\langle nlm|\mathbf{r}|n'l'm'\rangle = -e \int d^3r \psi_{nlm}^*(\mathbf{r}) \mathbf{r} \psi_{n'l'm'}(\mathbf{r}). \quad (3.4)$$

The atom can absorb a photon on the  $nl \rightarrow n'l'$  transition if this matrix element is nonzero.

- (1) Show that these matrix elements are zero if  $l$  and  $l'$  have the same parity.
- (2) Compute the matrix elements between the  $1s$  and  $2p$  states.

**Hydrogen wave functions.** They are of the form (Messiah, vol. 1, appendix B)

$$\psi_{nlm}(r, \theta, \varphi) = a_0^{-3/2} R_{nl}(r) Y_l^m(\theta, \varphi)$$

with the Bohr radius  $a_0$ , the radial wave function  $R_{nl}$  and the spherical harmonic (*Kugelflächenfunktion*)  $Y_l^m$ . For the states in question, the radial functions are given by

$$1s: \quad R_{10}(r) = 2e^{-r/a_0}$$

$$2s: \quad R_{20}(r) = \frac{\sqrt{2}}{2} \left(1 - \frac{r}{2a_0}\right) e^{-r/2a_0}$$

$$2p: \quad R_{21}(r) = \frac{\sqrt{6}r}{12a_0} e^{-r/2a_0}$$

The first few spherical harmonics are given by (Messiah's phase convention)

$$\text{s orbital: } Y_0^0(\theta, \varphi) = \frac{1}{\sqrt{4\pi}}$$

$$\text{p orbitals: } Y_1^0(\theta, \varphi) = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^{+1}(\theta, \varphi) = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\varphi}$$

$$Y_1^{-1}(\theta, \varphi) = \sqrt{\frac{3}{8\pi}} \sin \theta e^{-i\varphi}$$

For both questions, it is useful to compute first the integrals over the angles  $\theta$  and  $\varphi$ . For the radial integrals, you can use

$$\int_0^\infty dx x^n e^{-x} = \Gamma(n+1) = n!$$