

Problem 6.1 – P - and Q -function (10 points)

(a) In the lecture, we have computed the Q -function for a Fock (number) state:

$$Q_n(\alpha) = \frac{1}{\pi} |\langle \alpha | n \rangle|^2 = \exp(-|\alpha|^2) \frac{|\alpha|^{2n}}{\pi n!} \quad (6.1)$$

Using the relations derived in the lecture between the P - and Q -functions, show that the P -“function” for the Fock state is given by the highly singular expression [Sudarshan, *Phys Rev Lett* **10** (1963) 277; no guarantee for errors]

$$P_n(\alpha) = \frac{e^{|\alpha|^2}}{n!} \frac{\partial^n}{\partial \alpha^n} \frac{\partial^n}{\partial \alpha^{*n}} \delta^{(2)}(\alpha) \quad (6.2)$$

(b) Starting from Eq.(6.2), show that the average photon number in a Fock state $|n\rangle$ is given by $\dots n$ and that its variance vanishes.

Solution. The direct calculation with the help of a deconvolution is tricky because one ends up with an integral representation of derivatives of the δ -function, or with a binomial sum of derivatives:

$$P_n(\alpha) = \sum_{k=0}^n \binom{n}{k} \frac{1}{k!} \frac{\partial^k}{\partial \alpha^k} \frac{\partial^k}{\partial \alpha^{*k}} \delta^{(2)}(\alpha) \quad (6.3)$$

A more direct proof of Eq.(6.2) can be found in the book by Mandel & Wolf: observe that

$$\langle -\beta | \rho | \beta \rangle = \int d^2\alpha P(\alpha) \langle -\beta | \alpha \rangle \langle \alpha | \beta \rangle = e^{-|\beta|^2} \int d^2\alpha P(\alpha) e^{-|\alpha|^2} e^{\beta\alpha^* - \beta^*\alpha} \quad (6.4)$$

where the Fourier transform of the product $P(\alpha) e^{-|\alpha|^2}$ appears. The (formal) Fourier inversion then yields the representation

$$P(\alpha) = e^{|\alpha|^2} \int \frac{d^2\beta}{\pi^2} \langle -\beta | \rho | \beta \rangle e^{|\beta|^2} e^{-\beta\alpha^* + \beta^*\alpha} \quad (6.5)$$

which has, of course, very limited convergence properties as an ordinary function. Evaluating the matrix element $\langle -\beta | \rho | \beta \rangle$ for a Fock state is straightforward, and a formal calculation gives Eq.(6.2).

Problem 6.2 – Coherence of sunlight (10 points)

This is a problem to think about for the end of the year. The simplest model for the sunlight that reaches the earth is a multi-mode field that is made up from uncorrelated plane waves with some angular spread.

(a) Question: what is the spatial coherence function of sunlight, $g^{(1)}(\mathbf{r}, \mathbf{r}')$, after monochromatic filtering? You can use

$$\langle a_{\mathbf{k}}^\dagger a_{\mathbf{k}'} \rangle = \bar{n}(\omega_{\mathbf{k}}) \delta_{\mathbf{k}\mathbf{k}'} F(\hat{\mathbf{k}}) \quad (6.6)$$

where $F(\hat{\mathbf{k}})$ selects the relevant range of directions (of the order of 0.5° around the mean direction to the sun). To implement the frequency filter, include a δ -function fixing the mode frequencies to $\omega_{\mathbf{k}} = \omega$ under the k -summations/integrals.

(b) Determine the range of distances $\mathbf{r} - \mathbf{r}'$ where the coherence function $g^{(1)}(\mathbf{r}, \mathbf{r}')$ is not yet zero. (It decays to zero for large distances.) If you place two slits at \mathbf{r}, \mathbf{r}' where $|g^{(1)}(\mathbf{r}, \mathbf{r}')| \neq 0$, you should see Young's interference fringes (even without a 'coherence slit'). The corresponding coherence distance is surprisingly large.

(c) What value would you guess for the coherence distance, before embarking on the calculation?

Hint. For simplicity, take a screen perpendicular to the average direction towards the sun and consider only points on the screen. Ignore polarization, the refractive effect of the atmosphere, scattering by clouds etc.