

# Block vector & Rabi oscillations

Blatt # 3

real vector  
3 components

$s_{1,2,3}$  or  $u, v, w$

$$S_3 = \langle S_3 \rangle = \langle |e\rangle\langle e| - |g\rangle\langle g| \rangle = |c_e|^2 - |c_g|^2 \sim \sqrt{I} L$$

"Inversion"

$$s_1 + i s_2 = 2c_e^* c_g \quad (\text{oder } 2c_e c_g^*) = -\cos \Omega t$$

Projection of rotating vector  $\vec{S}(t)$

$$= \langle s_1 + i s_2 \rangle = 2 \langle |e\rangle\langle g| \rangle$$

$$= 2 c_e^* c_g \leftarrow \text{sensitive to superpositions}$$

$$\text{eg } 2c_e^* c_g = -i \sin(\Omega t)$$

on resonance (laser  $\omega_L = \omega_{A \text{ atom}}$ )

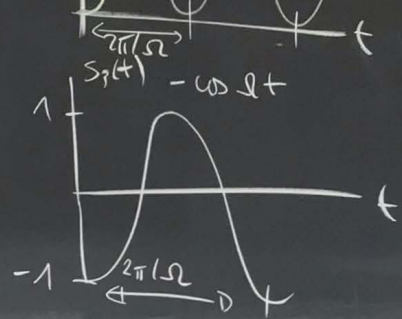
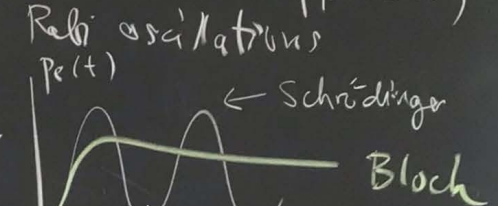
$$c_e(t) = i \sin(\Omega t/2)$$

$$c_g(t) = \cos(\Omega t/2)$$

$$\text{Rabi freq } \Omega = -\frac{d}{\hbar} \vec{E} \cdot \vec{L} \sim \sqrt{I} L$$

Python  
Bloch vector

rotating-wave = resonance (approx'n)



## Rabi oscillation

$$S_3(t) = -\cos \Omega t$$

$$S_2(t) = -\sin \Omega t$$

$$S_1(t) = 0$$

rotation ("precession") around  $S_1$ -axis ("effective magn. field")

$2\pi$  pulse (time  $\frac{2\pi}{\Omega}$ )  $|4\rangle \rightarrow -|4\rangle$

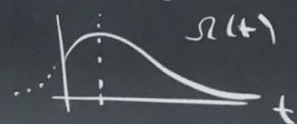
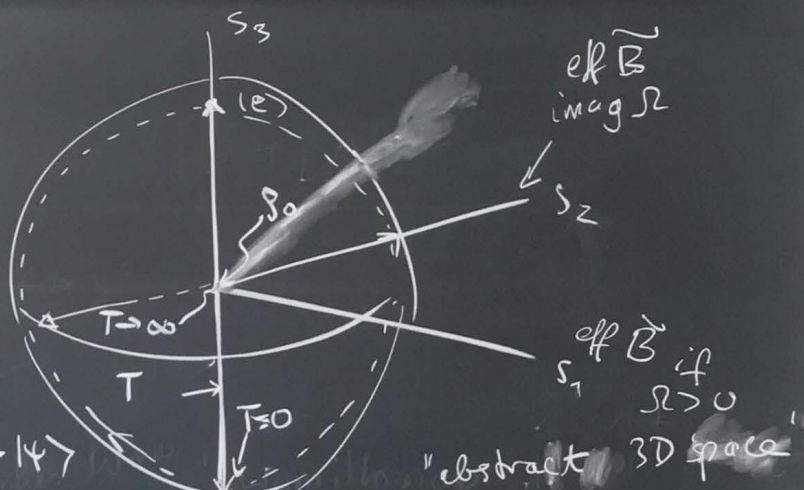
$\pi$  pulse  $(\pi/\Omega) |g\rangle \rightarrow i|e\rangle$

$\pi/2$  pulse  $(\pi/2\Omega) |g\rangle \rightarrow \frac{|g\rangle + i|e\rangle}{\sqrt{2}}$  (Rabi-flip)

(same weight for  $|g\rangle$  &  $|e\rangle$ )

Hadamard gate

Phase is fixed by complex  $\Omega$  adjust Rabi freq  $\Omega$  intera. time  $t$



pure states = sphere (surface) of Bloch "ball"

↑  
"wave functions"  $|ψ\rangle = c_0|e\rangle + c_1|g\rangle \leftarrow 3$  free parameters for both

$$S^2 = S_x^2 + S_y^2 + S_z^2 = (|c_0|^2 - |c_1|^2)^2 + 4|c_0 c_1|^2 = (|c_0|^2 + |c_1|^2)^2 = 1$$

(2 free parameters for pure states as Bloch vectors)  $4|c_0||c_1|^2$

outside - "forbidden" (pff #3.2) if  $|S| > 1$ , then probability  $\frac{1-|S|}{2} < 0$   
inner sphere - "mixed states"

generalises QM + class. statistics  
part of bigger system

↑  
eigenvalue of mixed state (density matrix)

density operator (matrix)

average  $p_1 \langle \psi_1 | A | \psi_1 \rangle + p_2 \langle \psi_2 | A | \psi_2 \rangle$

$= \text{tr} \left[ (p_1 |\psi_1\rangle\langle\psi_1| + p_2 |\psi_2\rangle\langle\psi_2|) A \right]$   $|\psi\rangle = \begin{cases} |g\rangle & \text{w/ } 99\% = p_1 \\ |e\rangle & \text{w/ } 1\% \text{ "error"} = p_2 \end{cases}$

↓ density matrix

$$\rho = 99\% |g\rangle\langle g| + 1\% |e\rangle\langle e| = \begin{pmatrix} 1\% & 0 \\ 0 & 99\% \end{pmatrix}$$

general basis

for  $\rho$ : eigenvectors

will differ from

incomplete knowledge preparation

↑ need not be orthogonal (!)

example

$$\vec{S} = 0$$

ex #3.2

$$\rho_0 = \frac{1}{2} + \frac{\vec{S} \cdot \vec{\sigma}}{2} = \begin{pmatrix} 1/2 & 0 \\ 0 & 1/2 \end{pmatrix} = \begin{cases} |e\rangle & 50\% \\ |g\rangle & 50\% \end{cases}$$

"thermal state" "completely mixed state"

$$= \begin{cases} \frac{|g\rangle + |e\rangle}{\sqrt{2}} & 50\% \\ \frac{|g\rangle - |e\rangle}{\sqrt{2}} & 50\% \end{cases}$$

$$\rho_T = \frac{e^{-\hbar\omega_A/kT} |e\rangle\langle e| + |g\rangle\langle g|}{Z}$$

$$Z = \text{partition function} = 1 + e^{-\hbar\omega_A/kT}$$

$$S_{1/2} = 0 \quad S_3 = k_B \ln \left( \frac{k_B T}{\hbar\omega_A} \right)$$

$$\hookrightarrow \text{tr}(\rho_T) = 1$$